Bessel Functions

1 A Special Class of Equations

Consider

$$(1 + R_M x^M) \frac{d^2 y}{dx^2} + \frac{1}{x} (P_0 + P_M x^M) \frac{dy}{dx} + \frac{1}{x^2} (Q_0 + Q_M x^M) y = 0$$
(1)

where, M is a positive integer. This equation is a special case of

$$\mathcal{L}y \equiv R(x)\frac{d^2y}{dx^2} + \frac{1}{x}P(x)\frac{dy}{dx} + \frac{1}{x^2}Q(x)y = 0$$
⁽²⁾

since $R(x) = 1 + R_M x^M$, $P(x) = P_0 + P_M x^M$ and $Q(x) = Q_0 + Q_M x^M$. Thus, x = 0 is a regular singular point. Let

$$y(x) = \sum_{k=0}^{\infty} a_k x^{k+s} \tag{3}$$

The Indicial Equation should be the same as before, i.e., for k = 0

$$f(s) = s^{2} + (P_{0} - 1)s + Q_{0} = 0$$
(4)

For $k = 1, 2, 3, \cdots$,

$$f(s+k)a_k + \sum_{i=1}^k g_i(s+k)a_{k-i} = 0$$
(5)

But

$$g_i(s) = R_i(s-i)(s-i-1) + P_i(s-i) + Q_i$$
(6)

for $i = 1, 2, 3, \dots$. Thus,

$$g_1(s) = g_2(s) = \dots = g_{M-1}(s) = g_{M+1}(s) = g_{M+2}(s) = \dots = 0$$
 (7)

since $R_i = P_i = Q_i = 0$ for $i = 1, 2, 3, \dots, M - 1, M + 1, \dots$, and

$$g_M(s) \neq 0 \tag{8}$$

due to the fact that $R_M \neq 0$, $P_M \neq 0$ or $Q_M \neq 0$.

When k = 1,

$$f(s+1)a_1 + g_1(s+1)a_0 = 0 \qquad \Rightarrow a_1 = 0 \tag{9}$$

When k = 2,

$$f(s+2)a_2 + \sum_{i=1}^{2} g_i(s+2)a_{2-i} = 0 \qquad \Rightarrow a_2 = 0 \tag{10}$$

Thus,

$$a_1 = a_2 = \dots = a_{M-1} = 0 \tag{11}$$

However, when k = M

$$f(s+M)a_M + \sum_{i=1}^{M} g_i(s+M)a_{M-i} = f(s+M)a_M + g_M(s+M)a_0 = 0$$
$$a_M = -\frac{g_M(s+M)}{f(s+M)}a_0$$
(12)

When k = M + 1,

$$f(s+M+1)a_{M+1} + g_M(s+M+1)a_1 = f(s+M+1)a_{M+1} = 0$$

$$a_{M+1} = 0$$
(13)

Similarly,

$$a_{M+1} = a_{M+2} = \dots = a_{2M-1} = 0 \tag{14}$$

When k = 2M,

$$f(s+2M)a_{2M} + g_M(s+2M)a_M = 0$$

$$a_{2M} = -\frac{g_M(s+2M)}{f(s+2M)}a_M = \frac{g_M(s+2M)}{f(s+2M)}\frac{g_M(s+M)}{f(s+M)}a_0$$
(15)

Thus,

$$a_{nM+1} = a_{nM+2} = \dots = a_{(n+1)M-1} = 0 \tag{16}$$

and

$$a_{(n+1)M} = -\frac{g_M(s + nM + M)}{f(s + nM + M)}a_{nM}$$
(17)

where $n = 0, 1, 2, \cdots$.

Therefore,

$$y(x) = \sum_{n=0}^{\infty} a_{nM} x^{nM+s} = \sum_{n=0}^{\infty} B_n x^{nM+s}$$
(18)

where $B_n = a_{nM}$.

The exceptional cases will be

- $s_1 = s_2$, or
- $s_1 s_2 = nM$, where *n* is a positive integer.

2 A Summary of Bessel Functions

The Bessel function satisfy

$$x^{2}\frac{d^{2}y}{dx^{2}} + x\frac{dy}{dx} + (x^{2} - p^{2})y = 0$$
(19)

where $p \ge 0$. Note that M = 2, $R_0 = 1$, $R_2 = 0$, $P_0 = 1$, $P_2 = 0$, $Q_0 = -p^2$ and $Q_2 = 1$.

• First Solution

$$J_p(x) = \sum_{k=0}^{\infty} \frac{(-1)^k (x/2)^{2k+p}}{\Gamma(k+p+1)k!}$$
(20)

- Second Solution
 - 1. $p \neq 0$ and p is not an integer

$$J_{-p}(x) = \sum_{k=0}^{\infty} \frac{(-1)^k (x/2)^{2k-p}}{\Gamma(k-p+1)k!}$$
(21)

2. p = 0

$$Y_0(x) = \frac{2}{\pi} \left[\left(\ln \frac{x}{2} + \gamma \right) J_0(x) - \sum_{k=1}^{\infty} (-1)^k \varphi(k) \frac{(x/2)^{2k}}{(k!)^2} \right]$$
(22)

where

$$\varphi(k) = \sum_{m=1}^{k} \frac{1}{m} \tag{23}$$

$$\gamma$$
(Euler constant) = $\lim_{k \to \infty} [\varphi(k) - \ln k] = 0.57721566 \cdots$ (24)

3. $p = \ell$ = a positive integer

$$Y_{\ell}(x) = \frac{2}{\pi} \left\{ \left(\ln \frac{x}{2} + \gamma \right) J_{\ell}(x) - \frac{1}{2} \sum_{k=0}^{\ell-1} \frac{(\ell - k - 1)! (x/2)^{2k-\ell}}{k!} - \frac{1}{2} \sum_{k=0}^{\infty} (-1)^{k} \left[\varphi(k) + \varphi(k+\ell) \right] \frac{(x/2)^{2k+\ell}}{k! (\ell+k)!} \right\}$$
(25)

3 Origins of Bessel Functions

Solutions of the differential equation

$$x^{2}\frac{d^{2}y}{dx^{2}} + x\frac{dy}{dx} + (x^{2} - p^{2})y = 0$$
(26)

or

$$x\frac{d}{dx}\left(x\frac{dy}{dx}\right) + (x^2 - p^2)y = 0 \tag{27}$$

are known as *Bessel functions of order* p, where p is real and non-negative.

Since M = 2, the solution is of the following form:

$$y(x) = \sum_{k=0}^{\infty} B_k x^{2k+s}$$
 (28)

After substituting this into equation (26), the resulting indicial equation, i.e., $f(s) = s^2 + (P_0 - 1)s + Q_0 = s^2 - p^2 = 0$, yields:

$$s_1 = p \qquad s_2 = -p \tag{29}$$

Thus, the exceptional cases occur only when

- $s_1 = s_2 = 0$, i.e., p = 0.
- $s_1 s_2 = 2p = 2\ell$, i.e., $p = \ell = a$ positive integer.

We can nonetheless always obtain the first solution for $s_1 = p$:

$$y_1(x) = B_0 \left[x^p + \sum_{k=1}^{\infty} \frac{(-1)^k x^{2k+p}}{2^{2k} k! (1+p)(2+p) \cdots (k+p)} \right]$$
(30)

According to the definition of gamma function, this solution can be written as

$$y_1(x) = B_0 \Gamma(1+p) \sum_{k=0}^{\infty} \frac{(-1)^k x^{2k+p}}{2^{2k} \Gamma(k+p+1)k!}$$
$$= 2^p \Gamma(1+p) B_0 \sum_{k=0}^{\infty} \frac{(-1)^k (x/2)^{2k+p}}{\Gamma(k+p+1)k!}$$
(31)

Here, let us define

$$J_p(x) = \sum_{k=0}^{\infty} \frac{(-1)^k (x/2)^{2k+p}}{\Gamma(k+p+1)k!}$$
(32)

or, if p is a non-negative integer, then

$$J_p(x) = \sum_{k=0}^{\infty} \frac{(-1)^k (x/2)^{2k+p}}{(k+p)!k!}$$

and $J_p(x)$ is known as the Bessel function of the first kind, of order p.

The second solution is presented in the sequel:

1. If $p \neq 0$ and is not an integer, a second solution is obtained by replacing p by -p in the first solution, i.e.,

$$J_{-p}(x) = \sum_{k=0}^{\infty} \frac{(-1)^k (x/2)^{2k-p}}{\Gamma(k-p+1)k!}$$
(33)

Thus the complete solution of equation (26) is

$$y(x) = c_1 J_p(x) + c_2 J_{-p}(x)$$
(34)

However, if p is an integer, it can be shown that

$$J_{-p}(x) = (-1)^p J_p(x)$$

Thus, $J_p(x)$ and $J_{-p}(x)$ are *linearly independent* in this situation. (Theorem 2, pp. 193, 194)

2. If p = 0, i.e., $s_1 = s_2 = 0$, then

$$y_2(x) = \left[\frac{\partial y(x,s)}{\partial s}\right]_{s=0}$$
(35)

The result of this derivation is

$$y_2(x) = B_0 \left[J_0(x) \ln x - \sum_{k=1}^{\infty} (-1)^k \varphi(k) \frac{(x/2)^{2k}}{(k!)^2} \right]$$
(36)

where

$$\varphi(k) = \sum_{m=1}^{k} \frac{1}{m}$$

Define

$$Y^{(0)}(x) = J_0(x) \ln x - \sum_{k=1}^{\infty} (-1)^k \varphi(k) \frac{(x/2)^{2k}}{(k!)^2}$$
(37)

Since equation (26) is linear, a linear combination of the above two solutions, i.e., $J_0(x)$ and $Y^{(0)}(x)$, is still a solution of the original differential equation. An alternative form of the second solution is thus more often used:

$$Y_0(x) = \frac{2}{\pi} \left[Y^{(0)} + (\gamma - \ln 2) J_0(x) \right]$$
$$= \frac{2}{\pi} \left[\left(\ln \frac{x}{2} + \gamma \right) J_0(x) + \sum_{k=1}^{\infty} (-1)^{k+1} \varphi(k) \frac{(x/2)^{2k}}{(k!)^2} \right]$$
(38)

where

$$\gamma = \lim_{k \to \infty} [\varphi(k) - \ln k] = 0.57721566490 \cdots$$

This standard particular solution is called *Bessel function of the second kind of* order zero or Neumann function of order zero.

This definition of the Bessel function of the 2nd kind is more convenient to use and thus is usually preferred, because of the fact that the behavior of the function $Y_0(x)$, for large value of x, is more comparable with the behavior of $J_0(x)$.

Thus, the complete solution in this case is

$$y(x) = c_1 J_0(x) + c_2 Y_0(x)$$
(39)

3. $p = \ell = a$ positive integer.

$$y_2(x) = \left\{ \frac{\partial}{\partial s} \left[(s+\ell)y(x,s) \right] \right\}_{s=-\ell}$$
(40)

The result

$$Y_{\ell}(x) = \frac{2}{\pi} \left\{ \left(\ln \frac{x}{2} + \gamma \right) J_{\ell}(x) - \frac{1}{2} \sum_{k=0}^{\ell-1} \frac{(\ell - k - 1)! (x/2)^{2k-\ell}}{k!} - \frac{1}{2} \sum_{k=0}^{\infty} (-1)^{k} \left[\varphi(k) + \varphi(k+\ell) \right] \frac{(x/2)^{2k+\ell}}{k! (\ell+k)!} \right\}$$
(41)

Notice that the second solution is defined differently, depending on whether the order p is an integer or not. To provide uniformity of formalism and numerical tabulation. it is desirable to adopt a form of the second solution that is *valid for all values of the order*. The standard second solution $Y_p(x)$ defined for all p is

$$Y_p(x) = \frac{1}{\sin p\pi} \left[J_p(x) \cos p\pi - J_{-p}(x) \right]$$
$$Y_{\ell}(x) = \lim_{p \to \ell} Y_p(x)$$

Asymptotic Behaviors of Bessel Functions

$$J_0(0) = 1 (42)$$

$$J_1(0) = J_2(0) = \dots = 0 \tag{43}$$

$$J_0(\infty) = J_1(\infty) = \dots = 0 \tag{44}$$

$$Y_0(0) = Y_1(0) = Y_2(0) = \dots = -\infty$$
(45)

$$Y_0(\infty) = Y_1(\infty) = \dots = 0 \tag{46}$$

Next, let $y = u/\sqrt{x}$, then substitute it into equation (26) to obtain

$$\frac{d^2u}{dx^2} + \left(1 - \frac{p^2 - 1/4}{x^2}\right)u = 0 \tag{47}$$

As $x \to \infty$,

$$1 \gg \frac{p^2 - 1/4}{x^2}$$

Thus,

$$\frac{d^2u}{dx^2} + u \simeq 0 \tag{48}$$

$$y \simeq \frac{1}{\sqrt{x}} (A\cos x + B\sin x) \tag{49}$$

It can be shown that, as $x \to \infty$,

$$J_p \sim \sqrt{\frac{2}{\pi x}} \cos\left(x - \alpha_p\right) \tag{50}$$

$$Y_p \sim \sqrt{\frac{2}{\pi x}} \sin\left(x - \alpha_p\right) \tag{51}$$

where,

$$\alpha_p = (2p+1)\frac{\pi}{4}$$

On the other hand, if p = 1/2, then equation (47) can be reduced to

$$\frac{d^2u}{dx^2} + u = 0\tag{52}$$

Thus, the complete solution is of the form

$$y = \frac{1}{\sqrt{x}} (A\cos x + B\sin x) \tag{53}$$

The standard solutions are

$$J_{1/2}(x) = \sqrt{\frac{2}{\pi x}} \sin x$$
 (54)

$$J_{-1/2}(x) = \sqrt{\frac{2}{\pi x}} \cos x$$
 (55)

4 Modified Bessel Functions

A slight variation of the standard Bessel equation (of order p) is the Bessel equation of order p with parameter α :

$$x^{2}\frac{d^{2}y}{dx^{2}} + x\frac{dy}{dx} + (\alpha^{2}x^{2} - p^{2})y = 0$$

This equation can be transformed into the standard form by substituting

 $t = \alpha x$

Thus,

$$t^{2}\frac{d^{2}y}{dt^{2}} + t\frac{dy}{dt} + (t^{2} - p^{2})y = 0$$

The solution of this equation is

$$y = c_1 J_p(t) + c_2 J_{-p}(t) = c_1 J_p(\alpha x) + c_2 J_{-p}(\alpha x)$$

or

$$y = c_1 J_p(t) + c_2 Y_p(t) = c_1 J_p(\alpha x) + c_2 Y_p(\alpha x)$$

Next, let us consider

$$x^{2}\frac{d^{2}y}{dx^{2}} + x\frac{dy}{dx} - (x^{2} + p^{2})y = 0$$
(56)

which is also an alternative form of the Bessel's equation. If we substitute t = ix into this equation, the resulting equation is

$$t^{2}\frac{d^{2}y}{dt^{2}} + t\frac{dy}{dt} + (t^{2} - p^{2})y = 0$$
(57)

which is the same as the standard Bessel's equation.

• If p is not zero or a positive integer, the general solution is

$$y(x) = c_1 J_p(ix) + c_2 J_{-p}(ix)$$
(58)

• Otherwise,

$$y(x) = c_1 J_{\ell}(ix) + c_2 Y_{\ell}(ix)$$
(59)

In the above equations, it is necessary to use

$$J_p(ix) = \sum_{k=0}^{\infty} \frac{(-1)^k}{\Gamma(k+p+1)k!} (\frac{ix}{2})^{2k+p} = i^p \sum_{k=0}^{\infty} \frac{1}{\Gamma(k+p+1)k!} (\frac{x}{2})^{2k+p}$$
(60)

Define

$$I_p(x) = \sum_{k=0}^{\infty} \frac{1}{\Gamma(k+p+1)k!} (\frac{x}{2})^{2k+p} = i^{-p} J_p(ix)$$
(61)

where $I_p(x)$ is referred to as the modified Bessel function of the first kind of order p. Thus, the complete solution for non-zero and non-integer p is

$$y(x) = c_1 I_p(x) + c_2 I_{-p}(x)$$
(62)

If $p = \ell = a$ non-negative integer, the second solution can be redefined as

$$K_{\ell}(x) = \frac{\pi}{2} i^{\ell+1} \left[J_{\ell}(ix) + iY_{\ell}(ix) \right]$$
(63)

where $K_{\ell}(x)$ is referred to as the modified Bessel function of the second kind of order ℓ .

Asymptotic Behavior of $I_p(x)$ and $K_p(x)$ as $x \to \infty$

$$I_p(x) \sim \frac{e^x}{\sqrt{2\pi x}} \to \infty$$
 (64)

$$K_p(x) \sim \frac{e^{-x}}{\sqrt{2x/\pi}} \to 0 \tag{65}$$

5 Properties of Bessel Functions

• For small values of x, i.e., as $x \to 0$, take the first term

$$J_p(x) \sim \frac{1}{2^p \Gamma(p+1)} x^p \tag{66}$$

$$J_{-p}(x) \sim \frac{2^p}{\Gamma(-p+1)} x^{-p} \qquad (p \neq \ell)$$
 (67)

$$Y_p(x) \sim -\frac{2^p(p-1)!}{\pi} x^{-p} \qquad (p \neq 0)$$
 (68)

$$Y_0(x) \sim \frac{2}{\pi} \ln x \tag{69}$$

$$I_p(x) \sim \frac{1}{2^p \Gamma(p+1)} x^p \tag{70}$$

$$I_{-p}(x) \sim \frac{2^p}{\Gamma(-p+1)} x^{-p} \qquad (p \neq \ell)$$
 (71)

$$K_p(x) \sim 2^{p-1}(p-1)! x^{-p} \qquad (p \neq 0)$$
 (72)

$$K_0(x) \sim -\ln x \tag{73}$$

Notice that only $J_p(x)$ and $I_p(x)$ are finite at x = 0 for $p \ge 0$.

• Differential properties:

$$\frac{d}{dx} [x^p y_p(\alpha x)] = \begin{cases} \alpha x^p y_{p-1}(\alpha x) & \text{for } y \equiv J, Y, I \\ -\alpha x^p y_{p-1}(\alpha x) & \text{for } y \equiv K \end{cases}$$

$$\frac{d}{dx} [x^{-p} y_p(\alpha x)] = \begin{cases} -\alpha x^{-p} y_{p+1}(\alpha x) & \text{for } y \equiv J, Y, K \\ \alpha x^{-p} y_{p+1}(\alpha x) & \text{for } y \equiv I \end{cases}$$
(74)

These formulas are established for J_p and Y_p by considering their series definitions, and for the remaining functions by considering their definitions in terms of J_p and Y_p .

From equation (74),

$$\frac{dy_p(\alpha x)}{dx} = \begin{cases} \alpha y_{p-1}(\alpha x) - (p/x)y_p(\alpha x) & \text{for } y \equiv J, Y, I \\ -\alpha y_{p-1}(\alpha x) - (p/x)y_p(\alpha x) & \text{for } y \equiv K \end{cases}$$
(76)

From equation (75),

$$\frac{dy_p(\alpha x)}{dx} = \begin{cases} -\alpha y_{p+1}(\alpha x) + (p/x)y_p(\alpha x) & \text{for } y \equiv J, Y, K\\ \alpha y_{p+1}(\alpha x) + (p/x)y_p(\alpha x) & \text{for } y \equiv I \end{cases}$$
(77)

- For J and Y,

By adding equations (76) and (77) and then dividing the result by 2, one can obtain

$$\frac{dy_p(\alpha x)}{dx} = \frac{\alpha}{2} \left[y_{p-1}(\alpha x) - y_{p+1}(\alpha x) \right]$$
(78)

By subtracting equation (77) from equation (76) and then dividing the result by 2, one can obtain

$$y_{p-1}(\alpha x) + y_{p+1}(\alpha x) = \frac{2p}{\alpha x} y_p(\alpha x)$$
(79)

- For I,

$$\frac{dI_p(\alpha x)}{dx} = \frac{\alpha}{2} \left[I_{p-1}(\alpha x) + I_{p+1}(\alpha x) \right]$$
(80)

$$I_{p-1}(\alpha x) - I_{p+1}(\alpha x) = \frac{2p}{\alpha x} I_p(\alpha x)$$
(81)

- For K,

$$\frac{dK_p(\alpha x)}{dx} = -\frac{\alpha}{2} \left[K_{p-1}(\alpha x) + K_{p+1}(\alpha x) \right]$$
(82)

$$K_{p-1}(\alpha x) - K_{p+1}(\alpha x) = -\frac{2p}{\alpha x} K_p(\alpha x)$$
(83)

[Example]

Using the table of J_0 and J_1 to integrate

$$I = \int_{1}^{2} x^{-3} J_4(x) dx$$

[Solution]

From (75) with p = 3 and $\alpha = 1$ we obtain

$$\frac{d}{dx}\left[x^{-3}J_3(x)\right] = -x^{-3}J_4(x)$$

i.e.,

$$\int_{1}^{2} x^{-3} J_4(x) dx = -x^{-3} J_3(x) \Big|_{x=1}^{x=2}$$

By (79), i.e., $J_{p-1}(\alpha x) + J_{p+1}(\alpha x) = \frac{2p}{\alpha x} J_p(\alpha x)$, with p = 2 and $\alpha = 1$ we have

$$J_1(x) + J_3(x) = \frac{4}{x}J_2(x) \implies J_3(x) = \frac{4}{x}J_2(x) - J_1(x)$$

Again by (79), i.e., $J_{p-1}(\alpha x) + J_{p+1}(\alpha x) = \frac{2p}{\alpha x} J_p(\alpha x)$, with p = 1 and $\alpha = 1$,

$$J_0(x) + J_2(x) = \frac{2}{x} J_1(x) \qquad \Rightarrow \qquad J_2(x) = \frac{2}{x} J_1(x) - J_0(x)$$

Thus,

$$J_3(x) = \frac{4}{x} \left[\frac{2}{x} J_1(x) - J_0(x) \right] - J_1(x) = \left(\frac{8}{x^2} - 1 \right) J_1(x) - \frac{4}{x} J_0(x)$$

From Table A1,

$$J_{1}(2) = 0.5767; \quad J_{0}(2) = 0.2239 \quad \Rightarrow \quad J_{3}(2) = J_{1}(2) - 2J_{0}(2) = 0.1289$$
$$J_{1}(1) = 0.4401; \quad J_{0}(1) = 0.7652 \quad \Rightarrow \quad J_{3}(1) = 7J_{1}(1) - 4J_{0}(1) = 0.0199$$
$$J_{1}(1) = 0.0199$$

$$I = -\left\lfloor \frac{1}{8}J_3(2) - \frac{1}{1}J_3(1) \right\rfloor = 0.0038$$

Note that we obtained previously

$$J_{\frac{1}{2}} = \sqrt{\frac{2}{\pi x}} \sin x$$
$$J_{-\frac{1}{2}} = \sqrt{\frac{2}{\pi x}} \cos x$$

and also

$$J_{p-1}(\alpha x) + J_{p+1}(\alpha x) = \frac{2p}{\alpha x} J_p(\alpha x)$$

[Example]

Determine $J_{\frac{3}{2}}(x)$ and $J_{-\frac{3}{2}}(x)$

[Solution]

From (79), (54) and (55) and let p = 1/2 and $\alpha = 1$, we get

$$J_{\frac{3}{2}}(x) = \frac{1}{x}J_{\frac{1}{2}} - J_{-\frac{1}{2}}(x) = \sqrt{\frac{2}{\pi x}} \left(\frac{\sin x}{x} - \cos x\right)$$

From (79), (54) and (55), and let p = -1/2 and $\alpha = 1$, we can obtain

$$J_{-\frac{3}{2}}(x) = -\frac{1}{x}J_{-\frac{1}{2}} - J_{\frac{1}{2}}(x) = -\sqrt{\frac{2}{\pi x}}\left(\frac{\cos x}{x} + \sin x\right)$$

 $\Rightarrow J_{\pm\frac{1}{2}}(x), J_{\pm\frac{3}{2}}(x), J_{\pm\frac{5}{2}}(x), \cdots \text{ are elementary}!!!$

6 Differential Equations Satisfied by Bessel Functions

Given that the solution of

$$X^{2}\frac{d^{2}\mathcal{Y}}{dX^{2}} + X\frac{d\mathcal{Y}}{dX} + (X^{2} - p^{2})\mathcal{Y} = 0$$
(84)

can be written in the form

$$\mathcal{Y} = Z_p(X) = \begin{cases} c_1 J_p(X) + c_2 J_{-p}(X) & p \neq 0 \text{ and } p \neq \text{ integer} \\ c_1 J_p(X) + c_2 Y_p(X) & p = 0 \text{ or } p = \text{ integer} \end{cases}$$
(85)

Let

$$X = f(x)$$
 $\mathcal{Y} = \frac{y}{g(x)}$

Note that

$$\frac{d}{dX} = \frac{d}{dx}\frac{dx}{dX} = \frac{1}{f'(x)}\frac{d}{dx}$$
(86)

Substitution into equation (84) yields

$$f\frac{d}{dx}\left[\frac{1}{f'}\frac{d}{dx}\left(\frac{y}{g}\right)\right] + \frac{d}{dx}\left(\frac{y}{g}\right) + \frac{f'}{f}(f^2 - p^2)\frac{y}{g} = 0$$
(87)

The solution of this equation is

$$y = g(x)Z_p[f(x)] \tag{88}$$

In particular, if we select

$$f(x) = \frac{\sqrt{d}}{s} x^s \tag{89}$$

$$g(x) = x^{(1-a)/2} \exp\left[-\frac{b}{r}x^r\right]$$
(90)

$$p = \frac{1}{s} \sqrt{\left(\frac{1-a}{2}\right)^2 - c} \tag{91}$$

where a, b, c, d, r and s are constants. In addition, it is required that

$$1 - a^2 \ge 4c$$
$$d \ne 0 \qquad r \ne 0 \qquad s \ne 0$$

Then equation (87) can be transformed to

$$x^{2}\frac{dy^{2}}{dx^{2}} + x(a+2bx^{r})\frac{dy}{dx} + \left[c+dx^{2s}-b(1-a-r)x^{r}+b^{2}x^{2r}\right]y = 0$$
(92)

The solution of this equation is

$$y = g(x)Z_p[f(x)] = x^{\frac{1-a}{2}} \exp\left[-\frac{b}{r}x^r\right]Z_p\left[\frac{\sqrt{d}}{s}x^s\right]$$
(93)

Thus, if it is possible to identify a particular second order differential equation with equation (92) by suitably choosing the constants a, b, c, d, r and s, the solution is immediately given in terms of Bessel function of order p.

[Example] $x^2y'' + xy' + (\lambda^2x^2 - p^2)y = 0$ $(\lambda \neq 0)$ Since $a = 1, b = 0, c = -p^2, d = \lambda^2, r = r$ and s = 1, then

$$p = \frac{1}{1} \left[\left(\frac{1-1}{2} \right)^2 - (-p^2) \right]^{0.5} = p$$

Thus,

$$g(x) = 1$$
$$f(x) = \lambda x$$
$$y = Z_p(\lambda x)$$

[Example] $x^2y'' + xy' - (\lambda^2 x^2 + p^2)y = 0$ $(\lambda \neq 0)$ Since $a = 1, b = 0, c = -p^2, d = -\lambda^2, r = r$ and s = 1, then

$$p = \frac{1}{1} \left[\left(\frac{1-1}{2} \right)^2 - (-p^2) \right]^{0.5} = p$$

Thus,

$$g(x) = 1$$
$$f(x) = i\lambda x$$
$$y = Z_p(i\lambda x)$$

[Exercise] $y'' + 3x^5y = 0$

Ans:

$$y = \sqrt{x} Z_{\frac{1}{7}} \left(\frac{2\sqrt{3}}{7} x^{\frac{7}{2}} \right)$$

[Exercise] $y'' + 5x^4y = 0$ Ans:

$$y = \sqrt{x} Z_{\frac{1}{6}} \left(\frac{\sqrt{5}}{3} x^3 \right)$$